Study of $p \bar{p} ightarrow \chi_{c1,2} ightarrow J/\psi \gamma$

Elisa Fioravanti

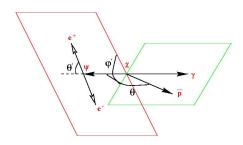
INFN Ferrara

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Study of $p\bar{p} \rightarrow \chi_{c1,2} \rightarrow J/\psi \gamma$

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χ_{c1} and χ_{c2} angular distributions



$\overline{p}p \rightarrow \chi_1 \rightarrow J/\psi\gamma$

- Production amplitudes: $B_0 = 0$
- Decay Amplitudes: a₂
 a₂ = 0.002 ± 0.032 ± 0.004

* E835 Collaboration, Nucl. Phys. B 717, 34 (2005)

- θ is the polar angle of the J/ψ with respect to the antiproton in the $\overline{p}p$ center of mass system

- θ' is the polar angle of the positron in the J/ψ rest frame with respect to the J/ψ direction in the χ rest of mass system

- ϕ' is the azimuthal angle between the J/ψ decay plane and the $\chi_{\rm c}$ plane

$\overline{p}p \rightarrow \chi_2 \rightarrow J/\psi\gamma$

- Production amplitudes: B_0^2 $B_0^2 = 0.16^{+0.09}_{-0.10} \pm 0.01$
- Decay Amplitudes: a_2, a_3 $a_2 = -0.076^{+0.054}_{-0.050} \pm 0.009$ $a_3 = 0.020^{+0.055}_{-0.044} \pm 0.009$

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$$\overline{\pmb{p}} \pmb{p}
ightarrow \chi_{1,2}
ightarrow \pmb{J}/\psi \gamma
ightarrow \ell^+ \ell^- \gamma$$

Cross section

$$\begin{split} \sigma(\chi_{c1} \to J/\psi\gamma) &\sim 1.7 \text{ nbarn} \\ \sigma(\chi_{c2} \to J/\psi\gamma) &\sim 2 \text{ nbarn} \\ \text{E835 Collaboration, Nucl. Phys. B 717, 34 (2005)} \\ \text{Background: } \bar{p}p \to \pi^+\pi^-\pi^0\text{: } \sigma(\chi_{c2}) = 0.12 \text{ mb} \\ \text{CERN-HERA 70-03 (1970)} \end{split}$$

Fast Simulation

•
$$J/\psi \rightarrow e^+e^-; J/\psi \rightarrow \mu^+\mu^-$$

- PID for Electrons: 1 Electron Loose; 1 Electron Tight (as in the Physics Book)
- PID for Muons: 1 Muon Loose; 1 Muon Tight (as in the Physics Book)
- PID for Photons: Neutral
- Bremsstrahlung effect for the electrons
- MC Truth Match
- 10.000 events generated
- Decay model: $\chi_{c1,2} \rightarrow J/\psi\gamma$: Chic1toJpsiGam
- Decay model: $J/\psi \rightarrow \ell^+ \ell^-$: VLL

4C fit is performed and best χ_{c1} candidate in each event is selected by minimal χ^2

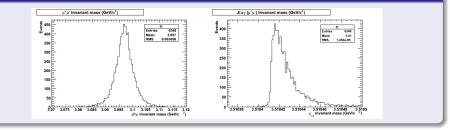
 $\overline{p}p \rightarrow \chi_{c1} \rightarrow J/\psi\gamma \rightarrow \ell^+\ell^-\gamma$

Invariant mass distributions

e*e' Invariant mass (GeV/c2) J/ψγ (e⁺e⁻) Invariant mass (GeV/c²) 350 3.51 3.096 Maar 300 1.125e-05 RMS 0.003901 250 200 150 100 50 3.07 3.075 3.08 3.085 3.09 3.095 3.1 0 3.51038 3.5104 3.51042 3.51044 3.51046 3.51048 3.105 3.11 3.115 3.12 3.5105 χ Invariant mass (GeV/c² J/w Invariant mass (GeV/c²)

 $J/\psi
ightarrow e^+e^-$

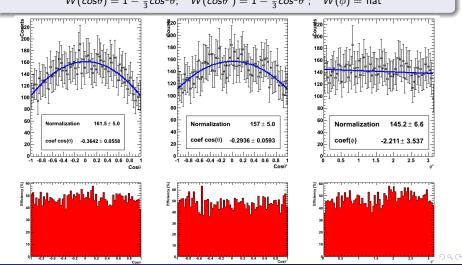
 $J/\psi \rightarrow \mu^+ \mu^-$



Efficiency $(J/\psi \rightarrow e^+e^-)$: 48.1%; Efficiency $(J/\psi \rightarrow \mu^+\mu^-)$: 63.5% = $-\infty$

Angular distributions for $J/\psi \rightarrow e^+e^-$

The angles distributions corrected with the efficiency, which is presented in the lower part. The angular distributions for the three angles can be approximately written as:

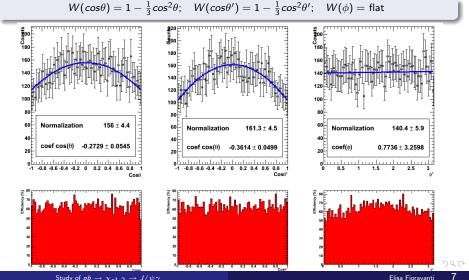


 $W(\cos\theta) = 1 - \frac{1}{2}\cos^2\theta; \quad W(\cos\theta') = 1 - \frac{1}{2}\cos^2\theta'; \quad W(\phi) = \text{flat}$

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Angular distributions for $J/\psi \rightarrow \mu^+\mu^-$

The angles distributions corrected with the efficiency, which is presented in the lower part. The angular distributions for the three angles can be approximately written as:



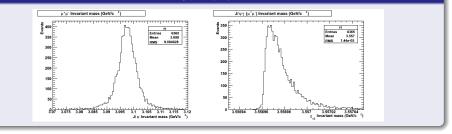
 $\overline{p}p \rightarrow \chi_{c2} \rightarrow J/\psi\gamma \rightarrow \ell^+\ell^-\gamma$

Invariant mass distributions

e'e' Invariant mass (GeV/c2) J/ψγ (e*e') Invariant mass (GeV/c2) 485 4855 3.557 Mean Mean 3.097 500 RMS 1.194e-05 RMS 0.004198 250 400 300 150 200 100 3.08 3.085 3.09 3.095 3.1 3.105 3.11 3.115 3.12 57 3.55702 3.55704 χ Invariant mass (GeV/c²) 3.07 3 075 3.55694 3.55696 3.55698 3.557 J/w Invariant mass (GeV/c²)

 $J/\psi
ightarrow e^+e^-$

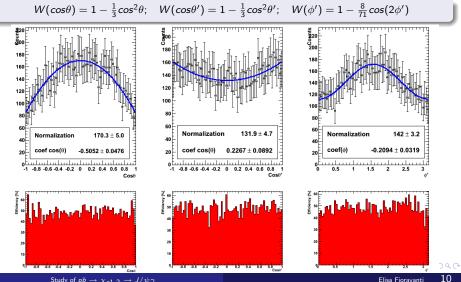
 $J/\psi \rightarrow \mu^+\mu^-$



Efficiency $(J/\psi \rightarrow e^+e^-)$: 48.6%; Efficiency $(J/\psi \rightarrow \mu^+\mu^-)$: 63.7%

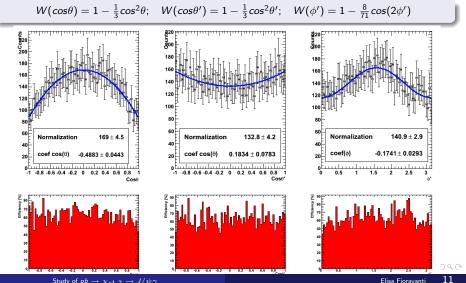
Angular distributions for $J/\psi \rightarrow e^+e^-$

The angles distributions corrected with the efficiency, which is presented in the lower part. The angular distributions for the three angles can be approximately written as:



Angular distributions for $J/\psi \rightarrow \mu^+\mu^-$

The angles distributions corrected with the efficiency, which is presented in the lower part. The angular distributions for the three angles can be approximately written as:



$\chi_{cj} ightarrow J/\psi\gamma$

- The angular distributions have been implemented in EvtGen
- The selection is in good shape.
- The reconstruction efficiency is:
 - ~48% for $J/\psi
 ightarrow e^+e^-$ (45% in the physics book)
 - \sim 63% for $J/\psi
 ightarrow \mu^+\mu^-$
- Next step: background studies

X8372) $\rightarrow J/\psi\pi\pi$

- I'm starting, I hope to show you the results the next week.

BACK-UP SLIDES

Image: A math a math

The measurement of the angular distributions in the radiative decays of the χ_c states provides the multipole structure of the radiative decay and the properties of the $\overline{c}c$ bound state.

$$\overline{\it p} \it p
ightarrow \chi_{\it c}
ightarrow {\it J}/\psi \gamma
ightarrow e^+ e^- \gamma$$

dominated by the dipole term E1.

M2 and E3 terms arise in the relativistic treatment of the interaction between the electromagnetic field and the quarkonium system. They contribute to the radiative width at the few percent level.

The angular distribution of the χ_1 and χ_2 are described by 4 indipendent parameters:

$a_2(\chi_{c1}), a_2(\chi_{c2}), B_0^2(\chi_{c2}), a_3(\chi_{c2})$

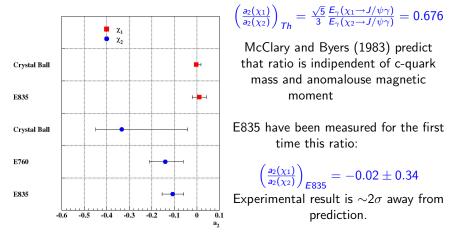
- The coupling between the set of χ states and $\overline{p}p$ is described by four indipendent helicity amplitudes:
 - χ_0 is formed only through the helicity 0 channel
 - χ_1 is formed only through the helicity 1 channel
 - χ_2 can couple to both
- The fractional electric octupole amplitude, $a_3 \approx E3/E1$, can contribute only to the χ_2 decays, and is predicted to vanish in the single quark radiation model if the J/ψ is pure S wave.
- For the fractional M2 amplitude a relativistic calculation yields:

$$a_2(\chi_{c1}) = -\frac{E_{\gamma}}{4m_c}(1+\kappa_c) = -0.065(1+\kappa_c)$$

$$a_2(\chi_{c2}) = -\frac{3}{\sqrt{5}}\frac{E_{\gamma}}{4m_c}(1+\kappa_c) = -0.096(1+\kappa_c)$$

where κ_c is the anomalous magnetic moment of the c-quark

χ_{c1} and χ_{c2} angular distributions



High statistics measurements of these angular distributions are needed to solve this question

E835 Reference "Ambrogiani et al. Physical Review D, Vol. 65, 05002"

For the χ_{c1} :

 $W(\cos\theta, \cos\theta', \phi') = \frac{1}{2}(1 - \cos^2\theta \cos 2\theta' - \sin\theta \cos\theta \sin\theta' \cos\theta' \cos\phi)$

For the χ_{c2} :

$$\begin{split} W(\cos\theta,\cos\theta',\phi') &= \frac{1}{10} [1 + \cos^2\theta + \cos^4\theta + 2\cos^2\theta\cos^2\theta' - 3\cos^4\theta\cos^2\theta' - \cos^2\theta\sin^2\theta\sin^2\theta'\cos^2\theta' + (\frac{1}{4}\sin2\theta - \sin2\theta\cos^2\theta)\sin2\theta'\cos\phi'] \end{split}$$

χ_{c1} and χ_{c2} angular distributions

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The angular distribution for the reactions (1) can be written as

$$W(\theta, \theta', \phi') = \sum_{i=1}^{N} K_i(B_{|\lambda(\bar{p}) - \lambda(p)|}, A_{|\lambda(\psi) - \lambda(\gamma)|})T_i(\theta, \theta', \phi')$$
(A1)

with N=5 at the χ_{c1} , and N=11 at the χ_{c2} . Tables IV (for the χ_{c1}) and V (for the χ_{c2}) give the full expressions (from Ref. [10]) for the coefficients K_i and the functions T_i that appear in Eq. (A1). The parameter R is defined as

TARIE IV	Coefficients K	and functions	T at the y

i	$T_i(\theta, \theta', \phi')$	$K_i(A_0, A_1)$
1	1	$\frac{1}{2}$
2	$\cos^2 \theta$	$\frac{1}{2}(A_1^2 - A_0^2)$
3	$\cos^2 \theta'$	$\frac{1}{2}(A_0^2-A_1^2)$
4	$\cos^2 \theta' \cos^2 \theta$	$-\frac{1}{2}$
5	$\sin 2\theta \sin 2\theta' \cos \phi'$	$-\frac{1}{4}A_0A_1$

$$\begin{pmatrix} A_0 = \frac{1}{\sqrt{2}} a_1 - \frac{1}{\sqrt{2}} a_2 \\ A_1 = \frac{1}{\sqrt{2}} a_1 + \frac{1}{\sqrt{2}} a_2 \end{pmatrix}_{J=1}$$

TABLE V. Coefficients K_i and functions T_i at the χ_{c2} .			
$T_i(\theta, \theta', \phi')$	$K_i(R,A_0,A_1,A_2)$		
1	$\frac{1}{2}(2A_0^2+3A_2^2-R(2A_0^2-4A_1^2+A_2^2))$		
$\cos^2 \theta$	$\frac{\frac{1}{8}(2A_0^2 + 3A_2^2 - R(2A_0^2 - 4A_1^2 + A_2^2))}{\frac{3}{4}(-2A_0^2 + 4A_1^2 - A_2^2 + R(4A_0^2 - 6A_1^2 + A_2^2))}$		
$\cos^4 \theta$	$\frac{1}{8}(6A_0^2 - 8A_1^2 + A_2^2)(3 - 5R)$		
$\cos^2 \theta'$			
$\cos^2 \theta' \cos^2 \theta$	$\frac{\frac{1}{8}(2A_0^2 + 3A_2^2 - R(2A_0^2 + 4A_1^2 + A_2^2))}{\frac{3}{4}(-2A_0^2 - 4A_1^2 - A_2^2 + R(4A_0^2 + 6A_1^2 + A_2^2))}$		
$\cos^2 \theta' \cos^4 \theta$	$\frac{1}{8}(6A_0^2+8A_1^2+A_2^2)(3-5R)$		
$\sin^2\theta'\cos 2\phi'$	$\sqrt{\frac{6}{4}}(R-1)A_0A_2$		
$\cos^2\theta\sin^2\theta'\cos 2\phi'$	$\sqrt{\frac{6}{4}}(4-6R)A_0A_2$		
$\cos^4\theta\sin^2\theta'\cos2\phi'$	$\sqrt{\frac{6}{4}}(5R-3)A_0A_2$		
$\sin 2\theta \sin 2\theta' \cos \phi'$	$-\sqrt{\frac{3}{4}}\left(A_0A_1+\sqrt{\frac{3}{2}}A_1A_2-R\left(2A_0A_1+\sqrt{\frac{3}{2}}A_1A_2\right)\right)$		
$\cos^2\theta\sin2\theta\sin2\theta'\cos\phi'$	$-\frac{1}{4\sqrt{3}}(5R-3)\left(3A_0A_1+\sqrt{\frac{3}{2}}A_1A_2\right)$		

$$\begin{vmatrix} A_0 = \sqrt{\frac{1}{10}}a_1 + \sqrt{\frac{1}{2}}a_2 + \sqrt{\frac{6}{15}}a_3 \\ \\ A_1 = \sqrt{\frac{3}{10}}a_1 + \sqrt{\frac{1}{6}}a_2 - \sqrt{\frac{8}{15}}a_3 \\ \\ A_2 = \sqrt{\frac{6}{10}}a_1 - \sqrt{\frac{1}{3}}a_2 + \sqrt{\frac{1}{15}}a_3 \end{vmatrix}$$