The DIRC detectors of the PANDA experiment at FAIR

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Abstract

The PANDA experiment at the future Facility for Anti-proton and Ion Research (FAIR) aims at studying strong interaction with antiprotons in the $1-15 \,\mathrm{GeV/c}$ range. For the charged particle identification, in particular kaons, one foresees three imaging Cherenkov detectors. Two DIRC detectors will be located in the target spectrometer section of PANDA, one a barrel shape DIRC with bar radiators in the centre part, and a disc DIRC in the forward endcap part. For the latter, two readout concepts are investigated: measuring the photons' time-of-propagation in a multi wavelength band disc DIRC, or measuring angles in a focussing lightguide dispersion-correcting disc DIRC.

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Key words: Particle identification, Cherenkov counter, ring imaging PACS: 29.40.Ka

The PANDA¹ experiment [1] at the future FAIR 1 laboratory is a fixed target experiment scattering 10 2 cooled antiprotons circulating in the High Energy 11 3 Storage Ring (HESR) with momenta of $1-15 \,\mathrm{GeV/c}_{-12}$ 4 off an internal pellet or gas jet target at interac- 13 5 tion rates of up to $2 \cdot 10^7$ /s to perform high preci-6 sion experiments in the charmed quark range. The 15 7 target spectrometer section with a superconducting 16 8

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solenoid and most subdetectors housed inside a return yoke covers the acceptance except for a hole of $\vartheta = 10^{\circ}$ horizontal and $\vartheta = 5^{\circ}$ vertical in the forward direction opening towards a dipole providing additional bending power and the subdetectors of the forward spectrometer section.

Three imaging Cherenkov detectors are foreseen for the charged particle identification of the PANDA experiment, two of them are DIRCs². For parti-17

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 $^{^1\,}$ anti
Proton ANnihilation at DArmstadt

² Detector of Internally Reflected Cherenkov light

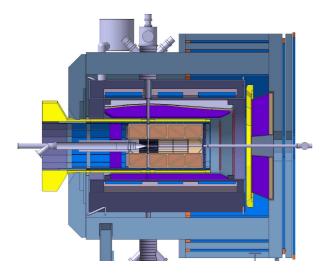


Fig. 1. PANDA target spectrometer. The Barrel and the Endcap DIRC detector positions are shown in light colour.

18 cle identification their information will be combined

¹⁹ with adjacent tracking and calorimetry detectors.

20 1. DIRC detectors

The PANDA target spectrometer (see Fig. 1) is 21 almost hermetically sealed to avoid solid angle gaps, 22 and to keep material volume low, there is little spare 23 space inside. The possibility of using thin radia-24 tors and placing the readout elements outside the 25 acceptance and potentially outside the magnet re-26 turn yoke favours DIRC designs to use as Cherenkov 27 imaging detectors for particle identification. 28

With the momentum ranges anticipated for the physics reactions in PANDA, Figure 2 showing a sample/benchmark reaction, the demands on the separation power increase towards forward angles.

Although the detectors cover different angu-33 lar ranges, each of the DIRC designs suggested 34 for PANDA nevertheless has to improve over cur-35 rently implemented DIRC designs, and hence needs 36 to address the effects of chromatic dispersion of 37 the Cherenkov light emission. For the time-of-38 propagation design selecting a narrow wavelength 39 band limits the group velocity time spread. For the 40 optical imaging design a prism element can largely 41 correct the dispersion angle spread. 42

The detectors are being located in a high radiation area, proton beam irradion up to 10Mrad [2] showing that amorphous fused silica is a radiation-hard radiator material. Several photon detector types are



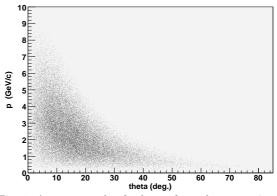


Fig. 2. Acceptance plot for kaons from the $p\bar{p} \rightarrow \psi_h \eta$ reaction at $p_{\bar{p}}=15 \,\text{GeV/c}$. The Barrel DIRC covers the range $\vartheta=22^{\circ}-140^{\circ}$, the Endcap DIRC the more forward range $\vartheta=5^{\circ}-22^{\circ}$.

 $B \approx 2$ Tesla [3] and high photon rates.

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The DIRC in the barrel part covers the angular range from $\vartheta = 22^{\circ}$ to 140°. The design is initially based on a scaled version the BaBar DIRC [4]. In order to improve performance, combining the time of arrival of the photons with their spatial image determines not only the particles velocity, but also the wavelength of the photons. Therefore dispersion correction at the lower and upper detection threshold is possible. See the separate paper [5] in these proceedings for details of the PANDA Barrel DIRC.

3. Endcap DIRC position

Two DIRC design options exist for the endcap part of the target spectrometer section. These differ in the photon readout design but both use an amorphous fused silica radiator disc. The endcap detector position covers forward angles up to $\vartheta = 22^{\circ}$ excluding an inner rectangular acceptance of $\vartheta_x = 10^{\circ}$ horizontal and $\vartheta_y = 5^{\circ}$ vertical half-angles.

In such a one-dimensional ³ DIRC type a photon is transported to the edge of a circular disc while preserving the angle information. Avoiding too much light scattering loss at the surface reflections requires locally (in the order of millimetres) a surface roughness not exceeding several nanometres RMS.

 $^{^{3}}$ Light is only reflected on surfaces of one spatial orientation, here the two disc surfaces both normal to the z axis.

The lower velocity threshold which is common to 74 both designs depends on the onset of total internal 75 reflection for a part of the photons emitted in the 76 Cherenkov cone.

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3.1. Time-of-Propagation disc DIRC 78

In the Multi-Chromatic Time-of-Propagation de-79 80 sign (see separate paper [6] for more detail) small detectors measure the arrival time of photons on the 81 disc rim requiring $\sigma_t=30-50$ ps single photon reso-82 lution. 83

For any given wavelength the disc edge is ef-84 fectively covered alternatingly with mirrors and 85 detectors. Only due to the resulting different light 86 path lenghts one can determine accurately enough 87 the start reference time, the time when the ini-88 tial charged particle enters the radiator and starts 89 creating the Cherenkov photons, as the stored an-90 tiproton beam in the HESR has no suitable time 91 structure to be used as an external time start. 92

As some of the light is reflected several times be-93 fore hitting a detector, the longer path lengths allow 94 a better relative time resolution. 95

The use of dicroic mirrors as colour filters allows 96 the use of multiple wavelength bands within the 97 same radiator (the current design suggesting two 98 bands) resulting in higher photon statistics. The 99 narrow wavelength bands minimise the dispersion 100 effects, and the quantum efficiency curve of the 101 photocathode material could be optimised for each 102 wavelength band individually. 103

3.2. Focusing Lightguide disc DIRC 104

In the Focusing Lightguide Dispersion-Correcting 105 design (Figures 3 and 4) when a photon arrives at 106 the edge of the circular or polygonal disc it enters 107 into one of about hundred optical elements on the 108 rim. Here the two-fold angle ambiguity (up-down) 121 109 is lifted, the chromatic dispersion corrected and the 122 110 photon focussed onto a readout plane. While the 123 111 optical element entered determines the ϕ coordinate, 124 112 measuring the position in the disperive direction on 125 113 the focal plane of the focusing lightguide yields the 126 114 θ coordinate. 115 127

Lithium fluoride (LiF) is UV transparent and has 128 116 particularly low-dispersion. Proton beam irradia- 129 117 tion of a test sample shows that radiation-produced 130 118 colour centres are confined to sufficiently small 131 119 wavelength ranges and only partially absorbing at 132 120

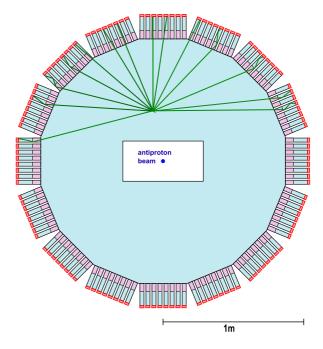


Fig. 3. Polygon disc with focussing lightguides attached to the rim used as optical readout components.

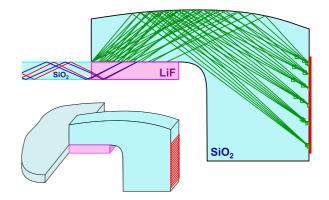


Fig. 4. Lightguide side view (inset 3D-visualisation) shown with a set of rays used for optimising the lightguide curvature. Reflections at the parallel front and back surfaces keep the light inside but do not affect the focussing properties.

the expected PANDA lifetime dose. Hence we think we can use LiF as a prism element (see Figure 4) to correct the Cherenkov radiation dispersion. The two boundary surfaces, with the radiator disc and the subsequent lightguide, make the chromatic dispersion correction angle-independent in first order.

As with the radiator, the light impinging on the inside of the lightguide's curved surface undergoes total internal reflection, hence no mirror coating is needed. The mirror makes the focussing also independent of the wavelength.

With the light staying within the dense optical

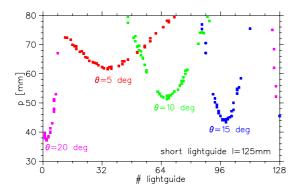


Fig. 5. Simulated photon hit pattern for four particles emitted at different angles θ and ϕ from the target vertex.

material of the lightguide most of the incoming light 133 phase space from the disc is be mapped onto the ¹⁶⁷ 134 focal plane with its one-coordinate readout. The fo-135 cussing surface with cylindrical shape of varying cur- 168 136 vature has been optimised to give an overall mini- 169 137 mum for the focus spot sizes of the different angles 170 138 on the focal plane, individual standard deviations 171 139 being well below 1 mm for the instrumented area. 172 140 173

141 3.3. Optical Simulations

Charged particles emitting cherenkov radiation
and photon propagation has been simulated as ¹⁷⁵
shown here for the Focussing Lightguide design, ¹⁷⁶
and detector resolution derived in analysing the ¹⁷⁷
photon hit patterns. ¹⁷⁸

The charged particle trajectory includes angular $^{179}\,$ 147 straggling, and the wavelength dependence of the 180 148 181 refractive indices is parametrised with Sellmeier co-149 efficients. Examples of idealisations: perfectly paral- $^{\ 182}$ 150 183 lel disc surfaces, no bulk light absorption and 100%151 reflectivity for total internal reflection (can be in-152 vestigated separately), no detector noise and light ¹⁸⁴ 153 from background particles. 154

The photon pattern analysis is seeded with particle vertex information smeared to the resolution 186 of the upstream tracking detectors. Differential re-

158 sponse vectors are computed, one tracking parame-

ter offset at a time, with high photon statistics using

the simulation code. Particle vertex parameters and
 velocity are then fitted simultaneously.

The detector resolving power is then derived from 190 the mean m and standard deviation σ of $\beta = v/c$ of 191 event samples for two different particle types A and 192 193

B of same momentum, yielding the sigma separation 194
 value defined as

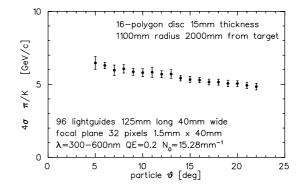


Fig. 6. Simulation-derived pion-kaon separation power for a focusing lightguide design with a 15 mm thick amorphous fused silica disc and 0.4 eV photon detection efficiency.

$$\sigma_{SEP} = \frac{|m_A - m_B|}{\sigma_\beta} = \frac{|m_A - m_B|}{(\sigma_A/2 + \sigma_B/2)}$$

For an Endcap DIRC detector with 96 lightguides and 3072 detector pixels that fits inside the target spectrometer return yoke Figure 6 shows the angledependent upper momentum limit being 5–6 GeV/c for 4σ pion-kaon separation within the acceptance $\vartheta = 5^{\circ}-22^{\circ}$.

4. Conclusions

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The high antiproton rates expected for the PANDA experiment require novel detectors. Current Research and Development addresses the point that the proposed detector components have to stand the harsh environment. We propose several DIRC detector designs for PID that fit into the limited available space of the target spectrometer, and with increased performance over currently running DIRC models meet the physics requirements.

5. Acknowledgements

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187

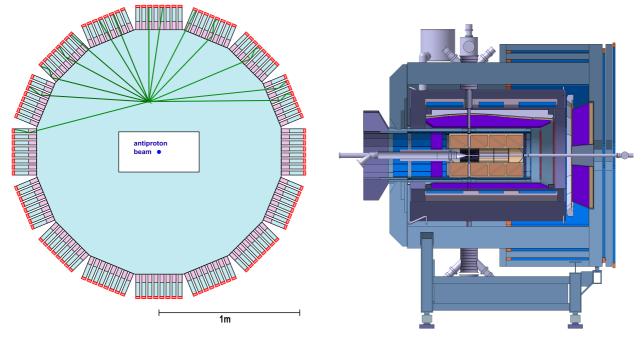


Fig. 7. disc 16 polygon.

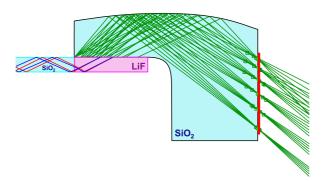


Fig. 8. lightguide side view with rays.

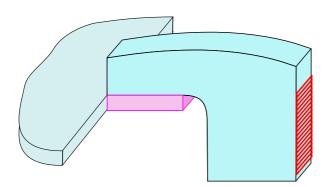


Fig. 9. Lightguide 3d visualisation.

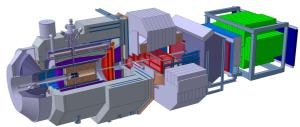


Fig. 10. PANDA target spectrometer

Fig. 11. PANDA spectrometer3d

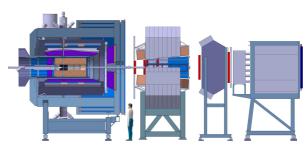


Fig. 12. PANDA spectrometer2d

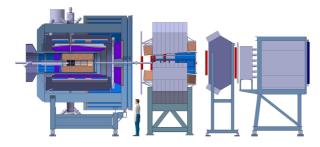


Fig. 13. PANDA spectrometer side view, target spectrometer left, further downstream the dipole in the middle with several forward spectrometer sections to the right.

The RICH detector in the forward spectrome-195 ter part will be of a standard aerogel radiator and 196 curved mirror design similar to the HERMES RICH 197 For the charged particle identification of the 198 PANDA experiment at FAIR, one foresees three 199 imaging Cherenkov detectors. An aerogel detector 200 of standard design will cover the forward spectrom-201 eter acceptance. 202

²⁰³ based on 15mm disc thickness, 96 lightguides ²⁰⁴ with 32 pixels each and 0.4eV (ϵ_{QE} =0.2 for λ =300-²⁰⁵ 600nm. The simulation results extend beyond the ²⁰⁶ actual angular coverage.

²⁰⁷ $p\bar{p} \rightarrow \psi_h \eta \; \vartheta = 22^{\circ} - 140^{\circ}$, the Endcap DIRC the ²⁰⁸ more forward range $\vartheta = 5^{\circ} = 22^{\circ}$.

Acceptance plot for kaons from the $p\bar{p} \rightarrow \psi_h \eta$ reaction at $p_{\bar{p}}=15 \text{ GeV/c}$. The Barrel DIRC covers the range $\vartheta=22^{\circ}-140^{\circ}$, the Endcap DIRC the more forward range $\vartheta=5^{\circ}=22^{\circ}$.

²¹³ ingredients: 3d simulation,

Sellmeier parametrisation of refractive indices wavelength dependence, analysis based on photon hit patterns, seeded with smeared Monte Carlo truth vertex equivalent to tracking information, resolution derived from event sample analysis distributions.

ideal geometry ideal light transmission, bulk, full
reflectivity (no Fresnel formula yet, photons unpolarised, hard cutoff for non-total internal reflection,

223 no background photons

224 explain 4 sigma